

Bregman divergences and error control via convex duality

Seminar - Numerical Mathematics

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Joint work with:



Alex Kaltenbach (TU Berlin)

- 📖 P.A. GAZCA–OROZCO, A. KALTENBACH *A Priori and A Posteriori Error Identities for Vectorial Problems via Convex Duality*. [ArXiv Preprint: 2602.04368](#), 2026.
- 📖 P.A. GAZCA–OROZCO *Bregman divergences and error control via convex duality*. In preparation.

2 The Laplace Equation

Let Ω be Lipschitz, $f \in L^2(\Omega)$

$$\begin{cases} -\Delta u = f & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega. \end{cases}$$

Find $u \in H_0^1(\Omega)$ s.t.:

$$\int_{\Omega} \nabla u \cdot \nabla v = \int_{\Omega} f v \quad \forall v \in H_0^1(\Omega).$$

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The solution $u \in H_0^1(\Omega)$ is a **minimiser** of the **convex energy** $I: H_0^1(\Omega) \rightarrow \mathbb{R}$

$$I(u) = \min_{v \in H_0^1(\Omega)} I(v)$$
$$I(v) := \frac{1}{2} \int_{\Omega} |\nabla v|^2 - \int_{\Omega} f v = \frac{1}{2} \|\nabla v\|_{\Omega}^2 - (f, v)_{\Omega}$$

2 The Laplace Equation (mixed I)

Introduce the **flux** $\mathbf{q} = \nabla u \in L^2(\Omega)^d$. One can then equivalently write:

Find $(\mathbf{q}, u) \in L^2(\Omega)^d \times H_0^1(\Omega)$ s.t.:

$$\left\{ \begin{array}{ll} -\operatorname{div} \mathbf{q} = f & \text{in } \Omega, \\ \mathbf{q} = \nabla u & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega. \end{array} \right. \quad \begin{array}{ll} -(\mathbf{q}, \mathbf{r})_{\Omega} + (\nabla u, \mathbf{r})_{\Omega} = 0 & \forall \mathbf{r} \in L^2(\Omega)^d \\ (\mathbf{q}, \nabla v)_{\Omega} = (f, v)_{\Omega} & \forall v \in H_0^1(\Omega). \end{array}$$

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The solution (\mathbf{q}, u) is a **saddle point** of the **Lagrangian**:

$$\begin{aligned} \max_{\mathbf{r} \in L^2(\Omega)} \mathcal{L}(u, \mathbf{r}) = \mathcal{L}(\mathbf{q}, u) &= \min_{v \in H_0^1(\Omega)} \mathcal{L}(\mathbf{q}, v) \\ \mathcal{L}(\mathbf{r}, v) &:= -\frac{1}{2} \|\mathbf{r}\|_\Omega^2 + (\mathbf{r}, \nabla v)_\Omega - (f, v)_\Omega \end{aligned}$$

2 Lagrangian duality

The **primal energy** can be written in terms of the Lagrangian:

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Lagrangian theory can help derive **dual (maximisation) problems**:

$$D(\mathbf{q}) = \sup_{\mathbf{r} \in L^2(\Omega)} D(\mathbf{r})$$
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In general there can be a **duality gap**:

$$D(\mathbf{r}) \leq D(\mathbf{q}) = \sup_{\mathbf{r} \in L^2(\Omega)} \inf_{v \in H_0^1(\Omega)} \mathcal{L}(\mathbf{r}, v) \leq \inf_{v \in H_0^1(\Omega)} \sup_{\mathbf{r} \in L^2(\Omega)} \mathcal{L}(\mathbf{r}, v) = I(u) \leq I(v)$$

2 Lagrangian duality

In general, the following are equivalent:

- ▶ (\mathbf{q}, u) is a saddle point of \mathcal{L}
- ▶ u is a solution of the **primal problem**, \mathbf{q} is a solution of the **dual problem**, there is **no duality gap**, i.e. $I(u) = D(\mathbf{q})$

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For Laplace, computing $\inf_{v \in H_0^1(\Omega)} \mathcal{L}(\mathbf{r}, v) = -\frac{1}{2} \|\mathbf{r}\|_{\Omega}^2 + (\mathbf{r}, \nabla v)_{\Omega} - (f, v)_{\Omega}$ leads to

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This implies that

$$\mathbf{r} \in H(\operatorname{div}; \Omega) := \{\mathbf{s} \in L^2(\Omega)^d \mid \operatorname{div} \mathbf{s} \in L^2(\Omega)\}$$
$$\operatorname{div} \mathbf{r} = -f$$

2 The Laplace equation (dual)

The **dual problem** corresponds to the **maximisation** problem:

$$D(\mathbf{q}) = \sup_{\mathbf{r} \in L^2(\Omega)} D(\mathbf{r})$$
$$D(\mathbf{r}) := -\frac{1}{2} \|\mathbf{r}\|_{\Omega}^2 - I_{\{-f\}}^{\Omega}(\operatorname{div} \mathbf{r})$$

The characteristic function $I_{\{-f\}}^{\Omega}$ is interpreted as:

$$I_{\{-f\}}^{\Omega}(\operatorname{div} \mathbf{r}) := \begin{cases} 0 & \text{if } \operatorname{div} \mathbf{r} = -f \\ +\infty & \text{otherwise.} \end{cases}$$

Note that we could maximise over $\mathbf{r} \in H(\operatorname{div}; \Omega)$.

2 The Laplace equation (mixed II)

Note that we could maximise over $\mathbf{r} \in H(\operatorname{div}; \Omega)$, and used the Lagrangian $\tilde{\mathcal{L}}: H(\operatorname{div}, \Omega) \times L^2(\Omega) \rightarrow \mathbb{R}$:

$$\tilde{\mathcal{L}}(\mathbf{r}, v) := -\frac{1}{2} \|\mathbf{r}\|_{\Omega}^2 - (\operatorname{div} \mathbf{r}, v)_{\Omega} - (f, v)_{\Omega}$$

This leads to the **optimality conditions**:

$$\begin{aligned} \partial_1 \tilde{\mathcal{L}}(\mathbf{q}, u) = 0 & \quad -(\mathbf{q}, \mathbf{r})_{\Omega} - (\operatorname{div} \mathbf{r}, u)_{\Omega} = 0 & \quad \forall \mathbf{r} \in H(\operatorname{div}; \Omega) \\ \partial_2 \tilde{\mathcal{L}}(\mathbf{q}, u) = 0 & \quad -(\operatorname{div} \mathbf{q}, v) = (f, v)_{\Omega} & \quad \forall v \in L^2(\Omega). \end{aligned}$$

The first condition implies that $\nabla u = \mathbf{q} \in L^2(\Omega)$, and so $u \in H_0^1(\Omega)$.

2 The Prager-Synge identity

The Prager-Synge (Hypercircle) Theorem [Prager, Synge; '47]

Let $v \in H_0^1(\Omega)$ and $\mathbf{r} \in H(\text{div}; \Omega)$ with $\text{div } \mathbf{r} = -f$ be arbitrary. Then

$$\frac{1}{2} \|\nabla u - \nabla v\|_{\Omega}^2 + \frac{1}{2} \|\mathbf{q} - \mathbf{r}\|_{\Omega}^2 = \frac{1}{2} \|\nabla v - \mathbf{r}\|_{\Omega}^2$$

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► The solution is characterised by three properties:

- > $u \in H_0^1(\Omega)$;
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- ▶ The RHS is computable \implies a posteriori estimation [Ainsworth, Oden, Braess, Schöberl, Bartels, Carstensen, Ern, Vohralík...].
- ▶ Geometric proof (Pythagoras) \implies difficult to generalise

2 Convex conjugates

The (convex) conjugate of a (proper) function $\varphi: \mathbb{R}^n \rightarrow \mathbb{R} \cup \{+\infty\}$ is the (proper, convex, lsc) function $\varphi^*: \mathbb{R}^n \rightarrow \mathbb{R} \cup \{+\infty\}$ defined by:

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For any proper convex lsc $\varphi: \mathbb{R}^n \rightarrow \mathbb{R} \cup \{+\infty\}$, the Fenchel–Young inequality holds:

$$\mathbf{a}^* \cdot \mathbf{a} \leq \varphi(\mathbf{a}) + \varphi^*(\mathbf{a}^*) \quad \forall \mathbf{a}, \mathbf{a}^* \in \mathbb{R}^n,$$

with equality being equivalent to the duality relations:

$$\mathbf{a}^* \cdot \mathbf{a} = \varphi(\mathbf{a}) + \varphi^*(\mathbf{a}^*) \quad \Leftrightarrow \quad \mathbf{a} = \nabla \varphi^*(\mathbf{a}^*) \quad \Leftrightarrow \quad \mathbf{a}^* = \nabla \varphi(\mathbf{a}).$$

(Or $\mathbf{a} \in \partial \varphi^*(\mathbf{a}^*) \Leftrightarrow \mathbf{a}^* \in \partial \varphi(\mathbf{a})$ for non-smooth φ .)

2 The generalised Prager-Synge identity

This can be generalised to **convex energies** on $W^{1,p}(\Omega)$:

$$I(v) := \int_{\Omega} [\varphi(\cdot, \nabla v) + \psi(\cdot, v)] \quad D(\mathbf{r}) = - \int_{\Omega} [\varphi^*(\cdot, \mathbf{r}) + \psi^*(\cdot, \operatorname{div} \mathbf{r})]$$

- 📖 S. REPIN *A posteriori estimates for partial differential equations*. Radon Series on Computational and Applied Mathematics 4, Walter de Gruyter GmbH & Co. KG, 2008.
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Define the total error:

$$\rho_{\text{tot}}^2(v, \mathbf{r}) := \rho_I^2(v) + \rho_{-D}^2(\mathbf{r}) := I(v) - I(u) + D(\mathbf{q}) - D(\mathbf{r})$$

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Thanks to strong duality, one concludes:

$$\rho_{\text{tot}}^2(v, \mathbf{r}) = I(v) - D(\mathbf{r}) =: \eta_{\text{gap}}^2(v, \mathbf{r})$$

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So the gap estimator is

$$\eta_{\text{gap}}^2(\mathbf{v}, \mathbf{r}) = \int_{\Omega} [\varphi(\cdot, \nabla \mathbf{v}) - \mathbf{r} \cdot \nabla \mathbf{v} + \varphi^*(\cdot, \mathbf{r})] + \int_{\Omega} [\psi(\cdot, \mathbf{v}) - \text{div } \mathbf{r} \mathbf{v} + \psi^*(\cdot, \text{div } \mathbf{r})]$$

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- ▶ Globally equivalent to **residual-type estimators**.

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- ▶ Globally equivalent to **residual-type estimators**.
- ▶ The integrands are **non-negative** and vanish if and only if the **optimality conditions** are satisfied.
- ▶ They can be **localised** \implies **adaptive mesh refinement**

2 The Prager-Synge identity (Laplace)

For Laplace, Taylor implies:

$$\rho_I^2(v) = I(v) - I(u) = \frac{1}{2} I''(u)[v - u, v - u] = \frac{1}{2} \|\nabla v - \nabla u\|_{\Omega}^2$$

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For the dual energy:

$$\rho_{-D}^2(\mathbf{r}) = D(\mathbf{q}) - D(\mathbf{r}) = \frac{1}{2} \|\mathbf{r}\|_{\Omega}^2 - \frac{1}{2} \|\mathbf{q}\|_{\Omega}^2$$

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So

$$\rho_{\text{tot}}^2(v, \mathbf{r}) = \frac{1}{2} \|\nabla v - \nabla u\|_{\Omega}^2 + \frac{1}{2} \|\mathbf{r} - \mathbf{q}\|_{\Omega}^2$$

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The duality gap estimator takes the form:

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$$\rho_{\text{tot}}^2(\mathbf{v}, \mathbf{r}) = \eta_{\text{gap}}^2(\mathbf{v}, \mathbf{r}) \iff \frac{1}{2} \|\nabla u - \nabla \mathbf{v}\|_{\Omega}^2 + \frac{1}{2} \|\mathbf{q} - \mathbf{r}\|_{\Omega}^2 = \frac{1}{2} \|\nabla \mathbf{v} - \mathbf{r}\|_{\Omega}^2$$

2 Local efficiency

The error estimator can be **localised**:

$$\frac{1}{2} \|\nabla v - \mathbf{r}\|_{\Omega}^2 = \sum_{K \in \mathcal{T}_h} \eta_K^2(\nabla v, \mathbf{r}) \quad \eta_K^2(\nabla v, \mathbf{r}) := \frac{1}{2} \int_K |\nabla v - \mathbf{r}|^2.$$

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For Laplace this is just the triangle inequality:

$$\begin{aligned} \eta_K^2(\nabla v, \mathbf{r}) &\leq \frac{1}{2} (\|\nabla v - \nabla u\|_K + \|\mathbf{q} - \mathbf{r}\|_K)^2 \\ &\leq \|\nabla v - \nabla u\|_K^2 + \|\mathbf{q} - \mathbf{r}\|_K^2 \\ &= 2\mathcal{E}_K \end{aligned}$$

2 What about more general energies?

The errors $I(v) - I(u)$ and $D(\mathbf{q}) - D(\mathbf{r})$ are **global quantities**, so it is not immediately obvious how to localise them.

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We will achieve this with the help of **Bregman divergences**:

$$\mathcal{D}_\varphi(\mathbf{a}, \mathbf{b}) := \varphi(\mathbf{a}) - \varphi(\mathbf{b}) - \nabla\varphi(\mathbf{b}) \cdot (\mathbf{a} - \mathbf{b}).$$

(or use a subgradient $\mathbf{r} \in \partial\varphi(\mathbf{b})$ in the non-smooth case.)

2 What about more general energies?

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This is a generalised notion of distance:

- ▶ **Not** necessarily symmetric $\mathcal{D}_\varphi(\mathbf{a}, \mathbf{b}) \neq \mathcal{D}_\varphi(\mathbf{b}, \mathbf{a})$.
- ▶ $\mathcal{D}_\varphi(\mathbf{a}, \mathbf{b}) \geq 0$ and $\mathcal{D}_\varphi(\mathbf{a}, \mathbf{a}) = 0$.
- ▶ If φ is strictly convex: $\mathcal{D}_\varphi(\mathbf{a}, \mathbf{b}) > 0$ if $\mathbf{a} \neq \mathbf{b}$.
- ▶ If $\mathbf{p} = \nabla\varphi(\mathbf{b})$, then $\mathcal{D}_\varphi(\mathbf{a}, \mathbf{b}) = \varphi(\mathbf{a}) + \varphi^*(\mathbf{p}) - \mathbf{p} \cdot \mathbf{a}$.
- ▶ For $\varphi(\mathbf{a}) = \frac{1}{2}|\mathbf{a}|^2$ we have $\mathcal{D}_\varphi(\mathbf{a}, \mathbf{b}) = \frac{1}{2}|\mathbf{a} - \mathbf{b}|^2$.

2 Updated error measures

We will replace our error measures:

$$I(v) - I(u) \quad \mapsto \quad \rho_I^2(v) := \int_{\Omega} \mathcal{D}_{\varphi}(\nabla v, \nabla u) + \int_{\Omega} \mathcal{D}_{\psi}(v, u)$$

$$D(\mathbf{q}) - D(\mathbf{r}) \quad \mapsto \quad \rho_{-D}^2(\mathbf{r}) := \int_{\Omega} \mathcal{D}_{\varphi^*}(\mathbf{r}, \mathbf{q}) + \int_{\Omega} \mathcal{D}_{\psi^*}(\operatorname{div} \mathbf{r}, \operatorname{div} \mathbf{q})$$

The integrands are now **pointwise non-negative** and **vanish if and only if $(v, \mathbf{r}) = (u, \mathbf{q})$** .

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Proposition [A. G; 2026]

Suppose strong duality $I(u) = D(\mathbf{q})$ is satisfied. Then the following **generalised Prager-Synge** identity holds:

$$\rho_I^2(v) + \rho_{-D}^2(\mathbf{r}) = \sum_{K \in \mathcal{T}_h} \eta_K^2(v, \mathbf{r})$$

$$\eta_K^2(v, \mathbf{r}) := \int_K [\varphi(\cdot, \nabla v) - \mathbf{r} \cdot \nabla v + \varphi^*(\cdot, \mathbf{r})] + \int_K [\psi(\cdot, v) - \operatorname{div} \mathbf{r} v + \psi^*(\cdot, \operatorname{div} \mathbf{r})],$$

Proposition [A. G.; 2026]

Suppose that the **Bregman–Young inequality** holds (here $\mathbf{r}_2 \in \partial\varphi(\mathbf{a}_2)$):

$$|(\mathbf{a}_1 - \mathbf{a}_2) \cdot (\mathbf{r}_1 - \mathbf{r}_2)| \leq c_{\text{BY}} (\mathcal{D}_\varphi(\mathbf{a}_1, \mathbf{a}_2) + \mathcal{D}_{\varphi^*}(\mathbf{r}_1, \mathbf{r}_2)).$$

Then the duality gap estimator is **locally efficient**:

$$\eta_{\varphi;K}^2(\nabla v, \mathbf{r}) \lesssim \int_K [\mathcal{D}_\varphi(\nabla v, \nabla u) + \mathcal{D}_{\varphi^*}(\mathbf{r}, \mathbf{q})],$$

where the constant **depends only on** c_{BY} .

The assumptions will be satisfied e.g. if $\varphi(\cdot, \mathbf{a}) = \widehat{\varphi}(|\mathbf{a}|)$, where $\widehat{\varphi}$ is a **uniformly convex N -function**.

2 Uniformly convex N -functions

A function $\widehat{\varphi}: \mathbb{R}_{\geq 0} \rightarrow \mathbb{R}_{\geq 0}$ is an N -function if [Diening, Ettwein, Růžička, Fornasier...]

- ▶ $\widehat{\varphi}$ is continuous and convex;
- ▶ There is a right-continuous, non-decreasing $\widehat{\varphi}': \mathbb{R}_{\geq 0} \rightarrow \mathbb{R}_{\geq 0}$ with

- > $\widehat{\varphi}(t) = \int_0^t \widehat{\varphi}'(s) ds$

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- > $\widehat{\varphi}'(s) > 0$ for $s > 0$

- > $\lim_{s \rightarrow \infty} \widehat{\varphi}'(s) = \infty$.

In particular $\lim_{t \rightarrow 0^+} \frac{\widehat{\varphi}(t)}{t} = 0$ and $\lim_{t \rightarrow +\infty} \frac{\widehat{\varphi}(t)}{t} = +\infty$.

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The N -function $\widehat{\varphi}$ is said to be **uniformly convex** if

$$C_{uc} \frac{\widehat{\varphi}'(s) - \widehat{\varphi}'(t)}{s - t} \leq \frac{\widehat{\varphi}'(s)}{s} \leq C_{uc} \frac{\widehat{\varphi}'(s) - \widehat{\varphi}'(t)}{s - t}$$

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Some examples:

$$\widehat{\varphi}(t) = \frac{t^p}{p} \quad \widehat{\varphi}(t) = t(\log(1+t))^p \quad p \in (1, \infty).$$

2 The φ -Laplacian

Consider the energy

$$I(v) := \int_{\Omega} \widehat{\varphi}(|\nabla v|) - \int_{\Omega} f v,$$

with uniformly convex $\widehat{\varphi}$. In this case there is a unique minimiser u , belonging to the Orlicz–Sobolev space:

$$W_0^{1,\varphi}(\Omega) := \left\{ v \in W_0^{1,1}(\Omega) \mid \nabla v \in L^\varphi(\Omega) \right\}$$
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The minimiser u satisfies

$$\int_{\Omega} \mathcal{A}(\nabla u) \cdot \nabla v = \int_{\Omega} fv \quad \forall v \in W_0^{1,\varphi}(\Omega).$$

with $\mathcal{A}(\nabla u) := \frac{\widehat{\varphi}'(|\nabla u|)}{|\nabla u|} \nabla u$ (With $\widehat{\varphi}(t) = \frac{t^p}{p}$ one has $\mathcal{A}(\nabla u) = |\nabla u|^{p-2} \nabla u$)

2 Uniformly convex N -functions

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- ▶ Local efficiency for the p -Laplacian in terms of the natural distance was known [Bartels, Kaltenbach; 2024], but the estimate contains **non-explicit constants**.
- ▶ Our error measure $\mathcal{D}_\varphi(\nabla v, \nabla u)$ is **equivalent** to the natural distance, and the estimate depends only on $c_{\text{uc}}, C_{\text{uc}}$.

2 Estimates for non-conforming functions

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- ▶ Define an **extended energy** on the broken Sobolev space $W^{1,1}(\mathcal{T}_h)$:

$$\begin{aligned} \bar{I}: W^{1,1}(\mathcal{T}_h) &\rightarrow \mathbb{R} \cup \{+\infty\} \\ \bar{I}(v) &:= \int_{\Omega} \varphi(\cdot, \mathfrak{G}_h v) + \int_{\Omega} \psi(\cdot, v), \end{aligned}$$

where \mathfrak{G}_h is a **discrete approximation of the gradient** (piecewise gradient, DG gradient...)

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- ▶ Similarly for the dual energy:

$$\bar{D}(\mathbf{r}) := - \int_{\Omega} \varphi^*(\cdot, \mathbf{r}) - \int_{\Omega} \psi^*(\cdot, \mathfrak{Div}_h \mathbf{r}),$$

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- ▶ **Solution:** Use Bregman divergences:

$$\begin{aligned}\bar{\rho}_{\bar{I}}(v) &:= \int_{\Omega} \mathcal{D}_{\varphi}(\mathfrak{G}_h v, \nabla u) + \int_{\Omega} \mathcal{D}_{\psi}(v, u) \\ \bar{\rho}_{-\bar{D}}(\mathbf{r}) &:= \int_{\Omega} \mathcal{D}_{\varphi^*}(\mathbf{r}, \mathbf{q}) + \int_{\Omega} \mathcal{D}_{\psi^*}(\mathfrak{D}iv_h \mathbf{r}, \operatorname{div} \mathbf{q})\end{aligned}$$

2 Estimates for non-conforming functions

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$$\bar{\rho}_I(v) + \bar{\rho}_{-D}(\mathbf{r}) = \sum_{K \in \mathcal{T}_h} \eta_K^2(v, \mathbf{r}) + (\mathbf{r} - \mathbf{q}, \mathfrak{G}_h v - \nabla u)_\Omega + (\mathfrak{D}iv_h \mathbf{r} - \operatorname{div} \mathbf{q}, \mathfrak{G}_h v - \nabla u)_\Omega$$

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- ▶ Hence, if \mathbf{r} is conforming \implies we can swap $u \mapsto s$ for an **arbitrary conforming** $s \in W^{1,1}(\Omega)$ in the remaining terms!
- ▶ We can then close the estimate with a **Young-type inequality** with ε .

Proposition [A. G.; 2026]

Suppose that the **Bregman–Young inequality with ε** holds (here $\mathbf{r}_2 \in \partial\varphi(\mathbf{a}_2)$):

$$|(\mathbf{a}_1 - \mathbf{a}_2) \cdot (\mathbf{r}_1 - \mathbf{r}_3)| \lesssim \varepsilon (\mathcal{D}_\varphi(\mathbf{a}_1, \mathbf{a}_2) + \mathcal{D}_{\varphi^*}(\mathbf{r}_1, \mathbf{r}_2)) + \mathcal{D}_\varphi(\mathbf{r}_1, \mathbf{r}_3).$$

(Similarly for ψ).

If \mathbf{r} is **div-conforming**, then we have the a posteriori bound:

$$\bar{\rho}_T(v) + \bar{\rho}_{-D}(\mathbf{r}) \lesssim \sum_{K \in \mathcal{T}_h} \eta_K^2(v, \mathbf{r}) + \inf_{s \in W_0^{1,1}(\Omega)} \left[\int_{\Omega} \mathcal{D}_\varphi(\mathfrak{G}_h v, \nabla s) + \int_{\Omega} \mathcal{D}_\psi(v, s) \right]$$

A similar statement holds if v is conforming.

► **Uniformly convex integrands** satisfy the assumptions.

2 Key properties

Crucial properties that we'd like to preserve at the **discrete level**:

- ▶ Poincaré inequality:

$$\|v\|_{\Omega} \lesssim \|\nabla v\|_{\Omega} \quad \forall v \in H_0^1(\Omega)$$

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- ▶ Orthogonality between gradients and divergence-free fields:

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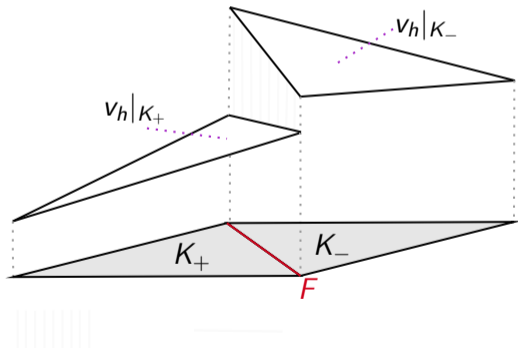
- ▶ Orthogonality $\implies H(\operatorname{div}; \Omega)$ regularity

$$\mathbf{r} \in L^2(\Omega)^d, \quad (\nabla v, \mathbf{r})_{\Omega} = 0 \quad \forall v \in H_0^1(\Omega) \quad \implies \quad \mathbf{r} \in H(\operatorname{div}; \Omega), \operatorname{div} \mathbf{r} = 0$$

2 Piecewise polynomials

On a **shape regular** mesh, denote the space of **piecewise polynomials of degree k** as:

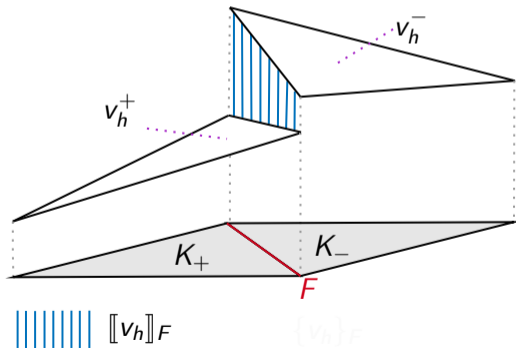
$$\mathbb{P}_k(\mathcal{T}_h) := \{v_h \in L^\infty(\Omega) \mid v_h|_K \in \mathbb{P}_k(K), \text{ for all } K \in \mathcal{T}_h\}$$



2 Jumps and averages

Define the **jump** of $v_h \in \mathbb{P}_k(\mathcal{T}_h)$ on a **facet** $F \in \mathcal{F}_h$ as:

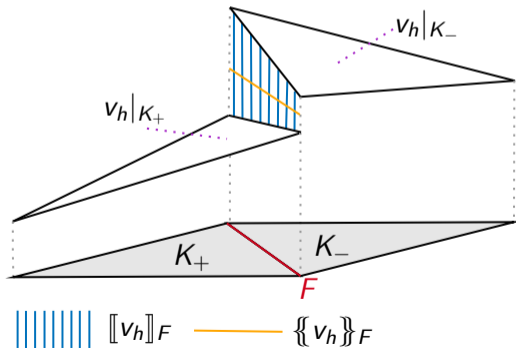
$$[[v_h]]_F := \begin{cases} v_h^+|_F - v_h^-|_F & \text{if } F \in \mathcal{F}_h^i, F = K_+ \cap K_- \\ v_h|_F & \text{if } F \in \mathcal{F}_h^\partial. \end{cases}$$



2 Jumps and averages

Define the **average** of $v_h \in \mathbb{P}_k(\mathcal{T}_h)$ on a **facet** $F \in \mathcal{F}_h$ as:

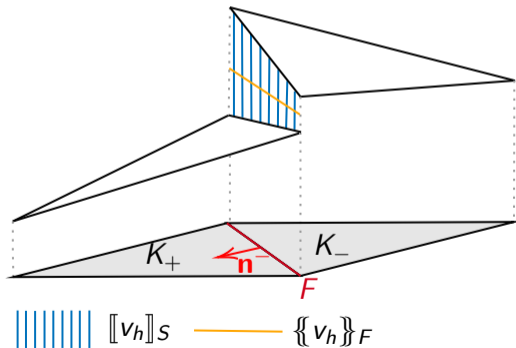
$$\{\{v_h\}\}_F := \begin{cases} \frac{1}{2} (v_h^+|_F + v_h^-|_F) & \text{if } F \in \mathcal{F}_h^i, F = K_+ \cap K_- \\ v_h|_F & \text{if } F \in \mathcal{F}_h^\partial. \end{cases}$$



2 Jumps and averages

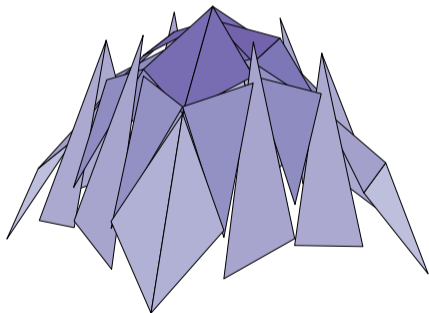
For a **vector-valued** function $\mathbf{r}_h \in \mathbb{P}_k(\mathcal{T}_h)^d$:

$$[[\mathbf{r}_h \cdot \mathbf{n}]]_F := \begin{cases} \mathbf{r}_h^+ \cdot \mathbf{n}^+ + \mathbf{r}_h^- \cdot \mathbf{n}^- & \text{if } F \in \mathcal{F}_h^i, F = K_+ \cap K_- \\ \mathbf{r}_h|_F & \text{if } F \in \mathcal{F}_h^\partial. \end{cases}$$

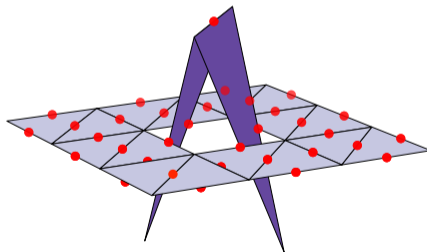


2 Crouzeix–Raviart element

$$V^h := \left\{ v_h \in \mathbb{P}_1(\mathcal{T}_h) \mid \int_F \llbracket v_h \rrbracket_F = \llbracket v_h \rrbracket_F(x_F) = 0 \text{ for all } F \in \mathcal{F}_h^i \right\}$$
$$V_D^h := \left\{ v_h \in V^h \mid \int_F v_h = v_h(x_F) = 0 \text{ for all } F \in \mathcal{F}_h^D \right\} \not\subset H_0^1(\Omega)$$



(a) Crouzeix–Raviart minimiser on $\Omega = (-1, 1)^2$ with $f \equiv 1$



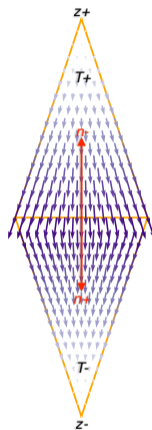
(b) Crouzeix–Raviart basis function

2 Raviart–Thomas element

$$\Sigma^h := \left\{ \mathbf{r}_h \in \mathbb{P}_1(\mathcal{T}_h)^d \mid \mathbf{r}_h \cdot \mathbf{n} \in \mathbb{P}_0(F) \text{ for all } F \in \mathcal{F}_h, \right. \\ \left. \llbracket \mathbf{r}_h \cdot \mathbf{n} \rrbracket_F = 0 \text{ for all } F \in \mathcal{F}_h^i, \right\}$$
$$\Sigma_N^h := \left\{ \mathbf{r}_h \in \Sigma^h \mid \mathbf{r}_h \cdot \mathbf{n} = 0 \text{ on } \Gamma_N \right\}$$

$$\Sigma^h \subset H(\text{div}; \Omega)$$

$$\text{div } \Sigma^h = \mathbb{P}_0(\mathcal{T}_h)$$



(a) RT basis function

2 Key properties at the discrete level

| 32

► Poincaré inequality (if $|\Gamma_D| > 0$):

$$\|v_h\|_{\Omega} \lesssim \|\nabla_h v_h\|_{\Omega} \quad \forall v_h \in V_D^h$$

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▶ Poincaré inequality (if $|\Gamma_D| > 0$): $\|v_h\|_{\Omega} \lesssim \|\nabla_h v_h\|_{\Omega} \quad \forall v_h \in V_D^h$

▶ Discrete integration by parts formula:

$$(\nabla_h v_h, \mathbf{r}_h)_{\Omega} + (v_h, \operatorname{div} \mathbf{r}_h)_{\Omega} = (v_h, \mathbf{r}_h \cdot \mathbf{n})_{\partial\Omega} \quad \forall v_h \in V^h, \mathbf{r}_h \in \Sigma^h.$$

▶ Orthogonality between discrete gradients and divergence-free fields:

$$(\nabla_h v_h, \mathbf{r}_h)_{\Omega} = 0 \quad \forall v \in V_D^h, \mathbf{r}_h \in \Sigma_N^h, \operatorname{div} \mathbf{r} = 0.$$

2 Key properties at the discrete level

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▶ Zero discrete divergence in $\mathbb{P}_0(\mathcal{T}_h)^d \implies$ Raviart–Thomas

$$\Pi_h \Sigma_N^h \perp_{L^2} \nabla_h(\ker \Pi_h|_{V_D^h})$$
$$(\mathbf{r}_h, \nabla_h v_h)_\Omega = 0 \quad \forall v_h \in V_D^h \text{ s.t. } \Pi_h v_h = 0 \implies \mathbf{r}_h = \Pi_h \tilde{\mathbf{r}}_h \text{ for some } \tilde{\mathbf{r}}_h \in \Sigma^h$$

 S. BARTELS, Z. WANG *Orthogonality relations of Crouzeix–Raviart and Raviart–Thomas finite element spaces. Numer. Math. 148:127–139, 2021*

2 Discrete strong duality

These properties allow the derivation of **discrete strong duality** principles for

$$I_h: V_D^h \rightarrow \mathbb{R} \cup \{+\infty\}$$
$$I_h(v_h) := \int_{\Omega} \varphi_h(\cdot, \nabla_h v_h) + \int_{\Omega} \psi_h(\cdot, \Pi_h v_h)$$

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The canonical discretisations would otherwise be:

- ▶ Mixed I: $(u, \mathbf{q}) \in H_D^1(\Omega) \times L^2(\Omega)^d \mapsto (u_h, \mathbf{q}_h) \in V_D^h \times \mathbb{P}_0(\mathcal{T}_h)^d$
- ▶ Mixed II: $(u, \mathbf{q}) \in L^2(\Omega) \times H_N(\operatorname{div}; \Omega)^d \mapsto (u_h, \mathbf{q}_h) \in \mathbb{P}_0(\mathcal{T}_h) \times \Sigma_N^h$

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In contrast, the duality approach delivers:

- ▶ Information about **gradients**: H^1 -convergence for u_h
- ▶ **Physically conforming** flux: $\mathbf{q} \in H(\operatorname{div}; \Omega)$

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It is possible to write optimality relations **pointwise**:

$$\Pi_h \mathbf{q}_h = D\varphi_h(\cdot, \nabla_h u_h) \quad \text{a.e. in } \Omega$$

$$\operatorname{div} \mathbf{q}_h = D\psi_h(\cdot, u_h) \quad \text{a.e. in } \Omega$$

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Crucially, one can obtain **generalised Marini formulas**:

$$\mathbf{q}_h = D\varphi_h(\cdot, \nabla_h u_h) + \frac{D\psi_h(\cdot, \Pi_h u_h)}{d} (\operatorname{id}_{\mathbb{R}^d} - \Pi_h \operatorname{id}_{\mathbb{R}^d})$$

2 A priori error analysis (Laplace)

For $f \in L^2(\Omega)$, find $u_h \in V_D^h$ s.t.

$$a_h(u_h, v_h) := (\nabla_h u_h, \nabla_h v_h) = (f, v_h)_\Omega$$

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Classical analysis [Berger, Scott, Strang; 1972]:

- ▶ Extend the bilinear form a_h to $H_0^1(\Omega) + V_D^h$
- ▶ With the help of an elliptic projection, prove the error bound

$$\|\nabla u - \nabla_h u_h\|_\Omega \leq \inf_{v_h \in V_D^h} \|\nabla u - \nabla v_h\|_\Omega + \sup_{v_h \in V_D^h} \frac{a_h(u - u_h, v_h)}{\|\nabla_h v_h\|_\Omega}$$

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- ▶ To handle the second term, note that

$$a_h(u - u_h, v_h) = - \sum_{F \in \mathcal{F}_h^i} \int_F \nabla u \cdot \mathbf{n}_F [[v]]_F$$

This requires $u \in H^{3/2+\varepsilon}(\Omega)$...

2 A priori error analysis (Laplace)

Alternative approach [Ern, Guermond; 2021]:

$$\|u - u_h\|_{\sharp} \leq c \inf_{v_h \in V_D^h} \|u - v_h\|_{\sharp} \quad \|v_h\|_{\sharp} := \sum_{K \in \mathcal{T}_h} \left(\|\nabla v\|_{\Omega}^2 + h_K \|\nabla v|_K \cdot \mathbf{n}\|_{\partial K}^2 \right)$$

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“Minimal regularity” estimates for $u \in H^1(\Omega)$ [Gudi; 2010]:

$$\|\nabla u - \nabla_h u_h\|_{\Omega} \leq c \left(\inf_{v_h \in V_h} \|\nabla u - \nabla_h v_h\|_{\Omega} + \text{osc}(f) \right)$$
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Truly minimal regularity ($u \in H_0^1(\Omega)$, $f \in H^{-1}(\Omega)$) estimates [Zanotti, Veerer; 2019]:

$$\|\nabla u - \nabla_h u_h\| \sim \inf_{v_h \in V_h} \|\nabla u - \nabla_h v_h\|_{\Omega} \quad \langle f, v_h \rangle_{\Omega} \mapsto \langle f, E_h v \rangle_{\Omega}$$

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What is c ?

2 A priori error analysis via duality (Laplace)

Convex duality easily delivers **constant-free** and (almost) **minimal-regularity** ($f \in L^2(\Omega)$) estimates [Bartels, Kaltenbach; 2024]:

$$\frac{1}{2} \|\nabla_h v_h - \nabla_h u_h\|_{\Omega}^2 + \frac{1}{2} \|\mathbf{r}_h - \mathbf{q}_h\|_{\Omega}^2 \leq \frac{1}{2} \|\mathbf{q}_h - \mathbf{r}_h\|_{\Omega}^2$$

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New idea: Mimic the a posteriori arguments (Prager–Synge) **at the discrete level**.

Estimate $(u, \mathbf{q}) - (v, \mathbf{r}) \longrightarrow$ Estimate $(u_h, \mathbf{q}_h) - (v_h, \mathbf{r}_h)$

\rightarrow recent work on the Signorini problem [Bartels, Gudi, Kaltenbach; 2024].

2 A priori error analysis via duality (Laplace)

$$I_h(v_h) := \frac{1}{2} \|\nabla_h v_h\|_{\Omega}^2 - (f_h, v_h)_{\Omega} \qquad D_h(\mathbf{r}_h) = -\frac{1}{2} \|\Pi_h \mathbf{r}_h\|_{\Omega}^2 - I_{\{-f_h\}}^{\Omega}(\operatorname{div} \mathbf{r}_h)$$

► **Step 1:** identify the strong convexity measures:

$$\begin{aligned} \rho_{I_h}^2(v_h) &:= I_h(v_h) - I_h(u_h) = \frac{1}{2} \|\nabla_h v_h - \nabla_h u_h\|_{\Omega}^2 \\ \rho_{-D_h}^2(\mathbf{r}_h) &:= D_h(\mathbf{q}_h) - D_h(\mathbf{r}_h) = \frac{1}{2} \|\Pi_h \mathbf{r}_h - \Pi_h \mathbf{q}_h\|_{\Omega}^2 \end{aligned}$$

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- **Step 2:** examine the discrete gap estimator:

$$\eta_{\text{gap},h}^2(v_h, \mathbf{r}_h) := I_h(v_h) - D_h(\mathbf{r}_h) = \frac{1}{2} \|\nabla_h v_h - \mathbf{r}_h\|_\Omega^2$$

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- **Step 3:** use the discrete strong duality $I_h(u_h) = D_h(\mathbf{q}_h)$:

$$\begin{aligned} \rho_{\text{tot},h}^2(v_h, \mathbf{r}_h) &= \eta_{\text{gap},h}^2(v_h, \mathbf{r}_h) \\ \frac{1}{2} \|\nabla_h u_h - \nabla_h v_h\|_{\Omega}^2 + \frac{1}{2} \|\Pi_h \mathbf{q}_h - \Pi_h \mathbf{r}_h\|_{\Omega}^2 &= \frac{1}{2} \|\nabla_h v_h - \Pi_h \mathbf{r}_h\|_{\Omega}^2 \end{aligned}$$

2 A priori error analysis via duality (Laplace)

Set $v_h = \mathcal{I}_h^{cr}(u)$ and $\mathbf{r}_h = \mathcal{I}_h^{rt}(\mathbf{q})$:

$$\|\nabla_h \mathcal{I}_h^{cr} u - \nabla_h u_h\|_{\Omega}^2 + \|\Pi_h \mathcal{I}_h^{rt} \mathbf{q} - \Pi_h \mathbf{q}_h\|_{\Omega}^2 = \|\nabla_h \mathcal{I}_h^{cr} u - \Pi_h \mathcal{I}_h^{rt} \mathbf{q}\|_{\Omega}^2$$

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Using the commuting property $\nabla_h \mathcal{I}_h^{cr} u = \Pi_h \nabla u = \Pi_h \mathbf{q}$, we obtain an **a priori error identity**:

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Similarly, we get **quasioptimality with explicit constants**:

$$\|\nabla u - \nabla_h u_h\|_{\Omega} + \|\mathbf{q} - \Pi_h \mathbf{q}_h\|_{\Omega} \sim \inf_{v_h \in V_D^h} \|\nabla u - \nabla_h v_h\|_{\Omega} + \inf_{\substack{\mathbf{r}_h \in \Sigma^h \\ \operatorname{div} \mathbf{r}_h = -f_h}} \|\mathbf{q} - \Pi_h \mathbf{r}_h\|_{\Omega}$$

2 A priori error analysis via duality (φ -Laplace)

Take a discretisation of the φ -Laplace problem:

$$I_h(v_h) := \int_{\Omega} \varphi(\nabla_h v_h) - (f_h, v_h)_{\Omega} - (F_h, \nabla_h v_h)_{\Omega} - (g_h, v_h)_{\Gamma_N},$$

$$D_h(\mathbf{r}_h) := - \int_{\Omega} \varphi^*(\Pi_h \mathbf{r}_h) - \chi_{\{-f_h\}}^{\Omega}(\operatorname{div}(\mathbf{r}_h - \mathbf{F}_h)) - \chi_{\{g_h\}}^{\Gamma_N}((\mathbf{r}_h - \mathbf{F}_h)\mathbf{n})$$

Here f_h, F_h, g_h are **piecewise constant approximations** of load terms $f, F, g \in L^{\varphi^*}$.

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Error bounds for Crouzeix–Raviart discretisations:

- ▶ [Yan, Liu, Carstensen...] Pioneering works with restrictive assumptions.
- ▶ [Kaltenbach; 2024], [Storn; 2025] *Medius analysis*: includes oscillation terms and shape-regularity-dependent constants.
- ▶ [Blechta, G-O, Kaltenbach, Růžička; 2026] Requires jump stabilisation and implementing a smoothing operator.
- ▶ [Diening, Hirn, Kreuzer, Zanotti; 2025] Focuses on Stokes (pressure-robustness), requires a smoothing operator and shape-regularity-dependent constants.

2 A priori error analysis via duality (φ -Laplace)

Proposition [A. G.; 2026]

The **discrete generalised Prager–Synge identity** holds for all $v_h \in V_D^h$ and $\mathbf{r}_h \in K_h$:

$$\int_{\Omega} \mathcal{D}_{\varphi}(\nabla_h u_h, \nabla u) + \int_{\Omega} \mathcal{D}_{\varphi^*}(\Pi_h \mathbf{q}_h, \mathbf{q}) = \int_{\Omega} [\varphi(\nabla_h v_h) - \Pi_h \mathbf{r}_h \cdot \nabla_h v_h + \varphi^*(\Pi_h \mathbf{r}_h)].$$

Furthermore, **quasioptimality** holds:

$$\int_{\Omega} \mathcal{D}_{\varphi}(\nabla_h u_h, \nabla u) + \int_{\Omega} \mathcal{D}_{\varphi^*}(\Pi_h \mathbf{q}_h, \mathbf{q}) \sim \inf_{v_h \in V_D^h} \int_{\Omega} \mathcal{D}_{\varphi}(\nabla_h v_h, \nabla u) + \inf_{\mathbf{r}_h \in K_h} \int_{\Omega} \mathcal{D}_{\varphi^*}(\Pi_h \mathbf{r}_h, \mathbf{q}),$$

where the constants **depend only on** c_{uc}, C_{uc} .

We get similar results for:

- ▶ Laplace (new for duality-based analysis):
 - > quasioptimality with known constants and minimal regularity
 - > Right-hand-side $f \in H^1(\Omega)$.
 - > Mixed inhomogeneous Dirichlet–Neumann BCs.
- ▶ Incompressible Stokes:
 - > A posteriori identity (known with geometric proof)
 - > Extended energy \Rightarrow Estimate for non-divergence-free velocities (known with geometric proof).
 - > Stress reconstruction formula (known with restrictive assumptions)
 - > Quasioptimality with known constants and minimal regularity (new).
- ▶ Linear elasticity:
 - > A posteriori identity (known with geometric proof)
 - > Extended energy \Rightarrow Estimate for non-symmetric stresses (known with geometric proof).
 - > Stress reconstruction formula (new)
 - > A priori estimate with known constants and minimal regularity (new).

